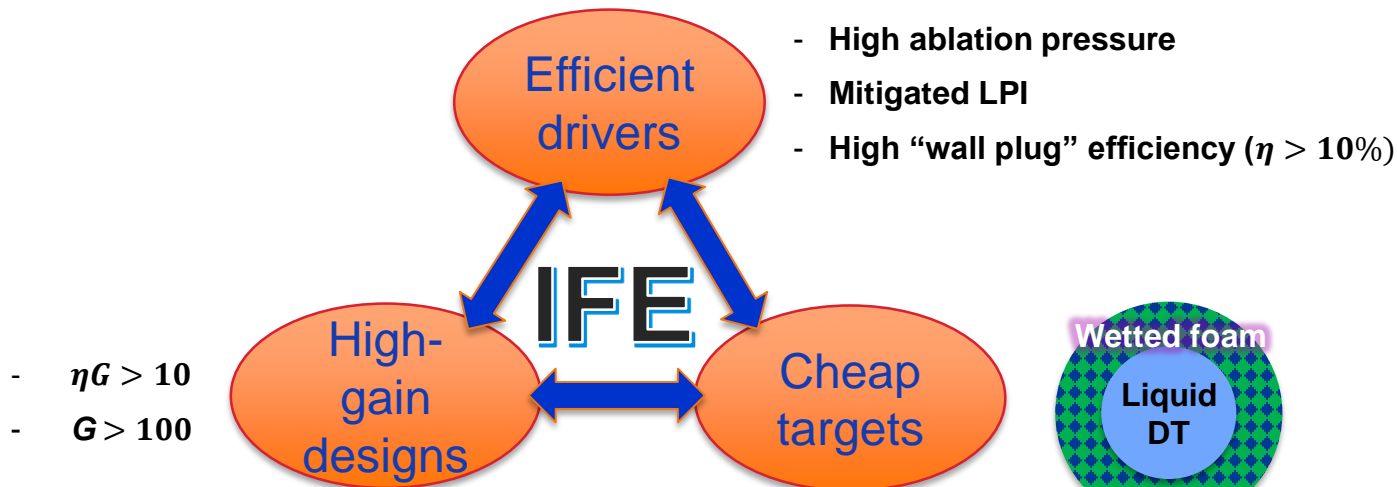


Inertial Fusion Energy Target Designs with Advanced Laser Technologies

IFE requirements (not unique!)

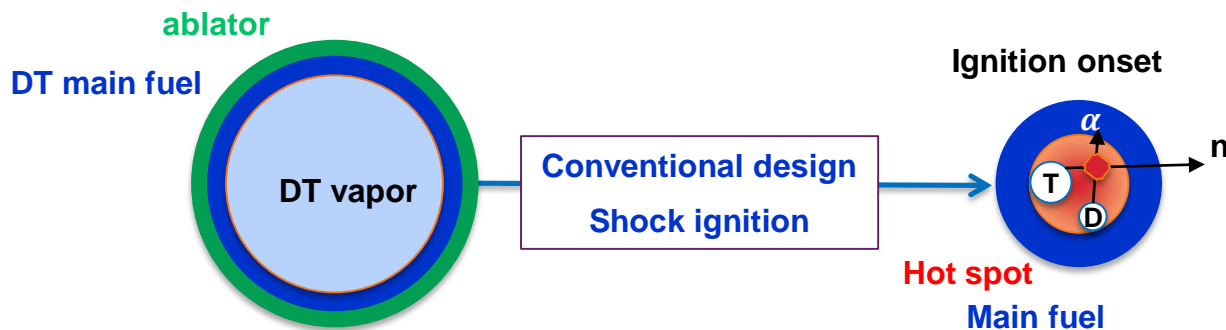


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Laboratory for Laser Energetics

IFE Science & Technology
Community Strategic Planning Workshop
22–24 February 2022


In the hot-spot ignition approach, the hot spot initiates a burn wave into the main fuel

Hot-spot ignition approach




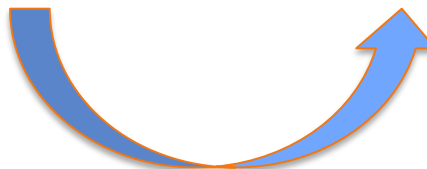
Compared to the current ICF implosions, the high-gain/IFE designs operate in somewhat different regions in the implosion parameter space

Current implosions

- 
- Maximize shell kinetic E_k and hot-spot E_{hs} energy
 - Maximize implosion velocity ($v_{imp} > 4 \times 10^7$ cm/s)
 - Fuel adiabat must be above a threshold value set by implosion stability
 - Penalty on convergence and ρR

High-gain designs

- 
- Fuel mass is maximized
 - Reduced implosion velocity ($v_{imp} < 3 \times 10^7$ cm/s)
 - Fuel must be kept close to Fermi degeneracy
 - Maximize convergence and ρR



How to bridge the gap?

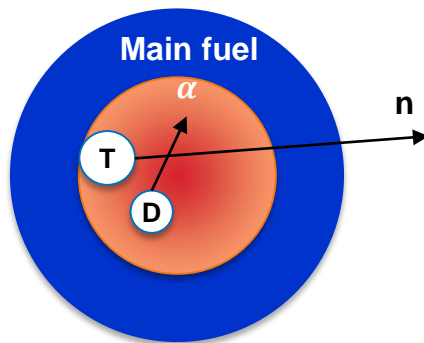
By increasing the drive pressure! IFE will require new drivers.

The drive pressure is key to the success of IFE

- Let's review IFE fundamentals

- Neutron yield: $Y_n = M_{\text{burn}} \epsilon_{\text{DT}}$, $\epsilon_{\text{DT}} = 2.75 \times 10^{11} \text{ J/g}$

$$M_{\text{burn}} = 3.7 \mu\text{g for } Y_n = 1 \text{ MJ}$$

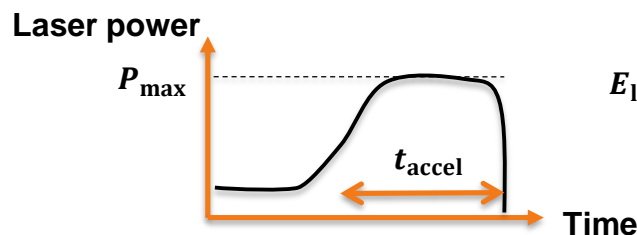
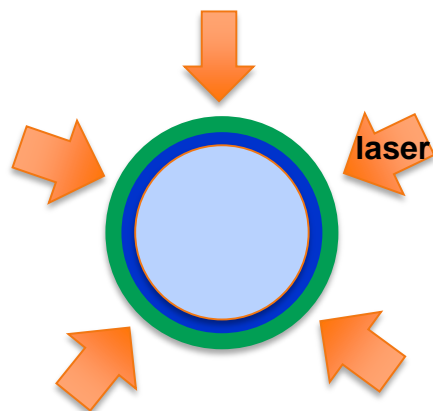


- If we want to get $Y_n = 100 \text{ MJ}$: $M_{\text{burn}} = 0.37 \text{ mg}$
- Burn fraction $f_{\text{burn}} \sim 1/4$: $M_{\text{DT}} = 1.5 \text{ mg}$
- We want to accelerate this mass to $v_{\text{imp}} = 3 \times 10^7 \text{ cm/s}$
- We will need $E_k = (M_{\text{DT}} v_{\text{imp}}^2)/2 = 70 \text{ kJ}$
- Current NIF: $E_k \sim 20 \text{ kJ}$
- Scaled LDD OMEGA implosions to NIF energies : $E_k = 60 \text{ to } 80 \text{ kJ}$

Energetically, LDD makes more sense for IFE but it must remove one obstacle first!

Accelerating the required amount of fuel sets the requirement on ablation pressure

Laser irradiation provides drive pressure p_a through mass ablation



$$E_{\text{laser}} \simeq P_{\text{max}} t_{\text{accel}}$$

$$M_{\text{DT}} \frac{v_{\text{imp}}}{t_{\text{accel}}} \simeq 4\pi R_{\text{target}}^2 p_a$$

$$M_{\text{DT}} \sim E_{\text{laser}} \frac{p_a}{v_{\text{imp}}} \underbrace{\left(\frac{4\pi R^2}{P_{\text{max}}} \right)}$$

Peak laser intensity I_{max}^{-1}

$$M_{\text{DT}} \simeq \frac{1}{5} \frac{E_{\text{laser}}}{I_{\text{max}}} \frac{p_a}{v_{\text{imp}}}$$

To accelerate $M_{\text{DT}} = 1.5 \text{ mg}$ to $v_{\text{imp}} = 3 \times 10^7 \text{ cm/s}$
at $E_{\text{laser}} \sim 1 \text{ MJ}$ requires $p_a > 220 \text{ Mbar}$

Target gain is a simple relation between ablation pressure, implosion velocity, and peak drive intensity

$$M_{\text{DT}} \simeq \frac{1}{5} \frac{E_{\text{laser}}}{I_{\text{max}}} \frac{p_a}{v_{\text{imp}}}$$

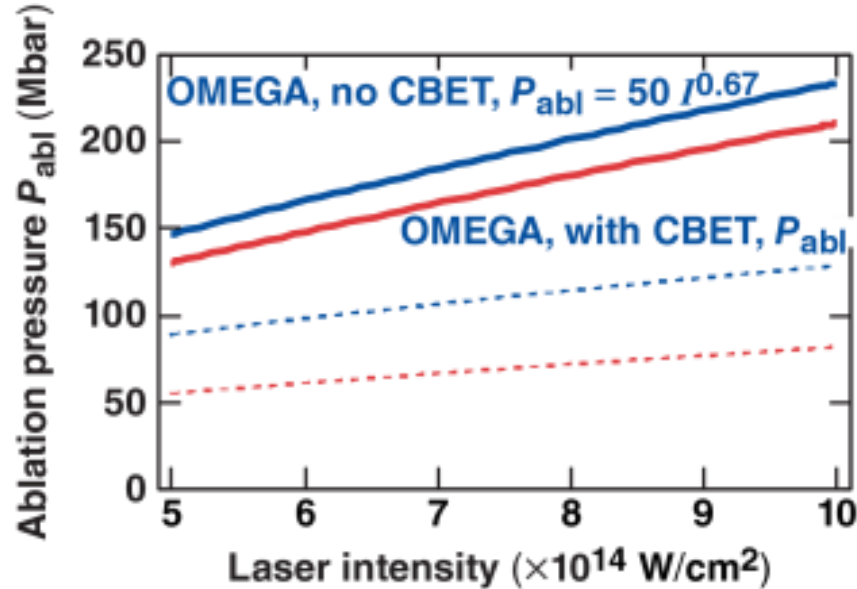
$$G = f_{\text{burn}} M_{\text{DT}} \epsilon_{\text{DT}} / E_L$$

$$G \simeq 180 \frac{f_{\text{burn}}}{I_{15}} \left(\frac{p_a}{100 \text{ Mbar}} \right) \left(\frac{v_{\text{imp}}}{3 \times 10^7 \text{ cm/s}} \right)^{-1}$$

For LDD implosions $p_a = p_a(I)$ depends on thermal conduction and LPI.

Ablation pressure depends on physics details

Direct Drive, $\lambda_L = 351 \text{ nm}$



$$G \approx 180 \frac{f_{\text{burn}}}{I_{15}} \left(\frac{p_a}{100 \text{ Mbar}} \right) \left(\frac{v_{\text{imp}}}{3 \times 10^7 \text{ cm/s}} \right)^{-1}$$

Bremsstrahlung absorption:

$$G \approx 95 I_{15}^{-0.31} \left(\frac{v_{\text{imp}}}{3 \times 10^7} \right)^{-1}$$

Bremsstrahlung + CBET losses:

$$G_{\text{CBET}} \approx 37 I_{15}^{-0.43} \left(\frac{v_{\text{imp}}}{3 \times 10^7} \right)^{-1}$$

$$I_{15} \sim 0.1, p_a \sim 20 \text{ Mbar}$$

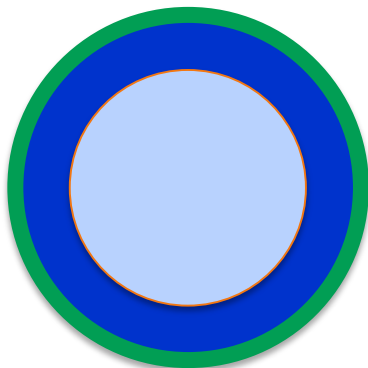
Maximizing fuel mass pushes the design to use lower drive intensities.

What's wrong with using lower drive intensities and pressures

I. Target Stability

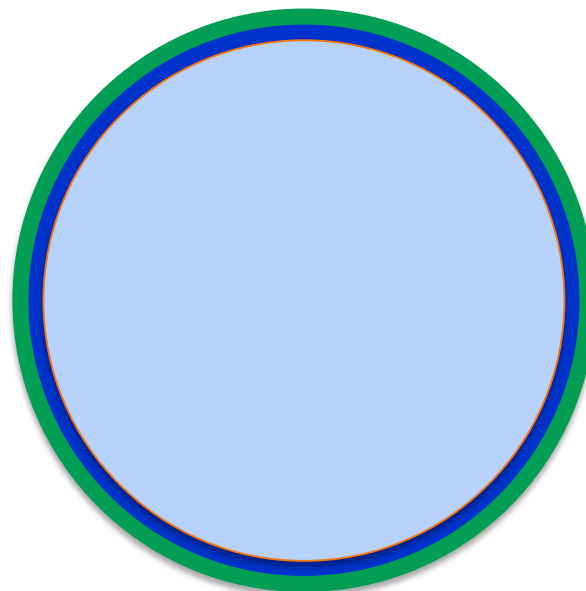
*severe constraints on target
quality and beam uniformity*

High drive pressure



$$IFAR \sim \frac{v_{\text{imp}}^2}{p_a^{\frac{2}{5}} \alpha^{\frac{3}{5}}}$$

Low drive pressure, longer acceleration distance



To accelerate fuel mass M_{DT} with lower drive pressure requires high-aspect-ratio shells.

What's wrong with using lower drive intensities and pressures

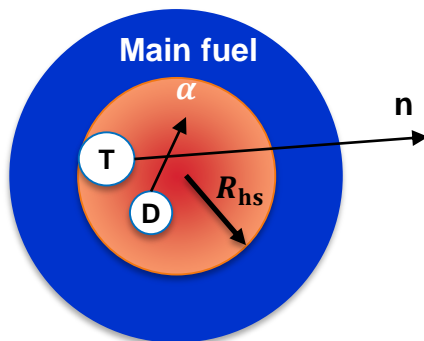
II. Creating an igniting hot spot is challenging

Lower drive pressure requires more laser energy for ignition

$$(\rho R)_{\text{hs}} T_{\text{hs}} = 0.3 \frac{\text{g}}{\text{cm}^2} \times 5 \text{ keV} \rightarrow$$

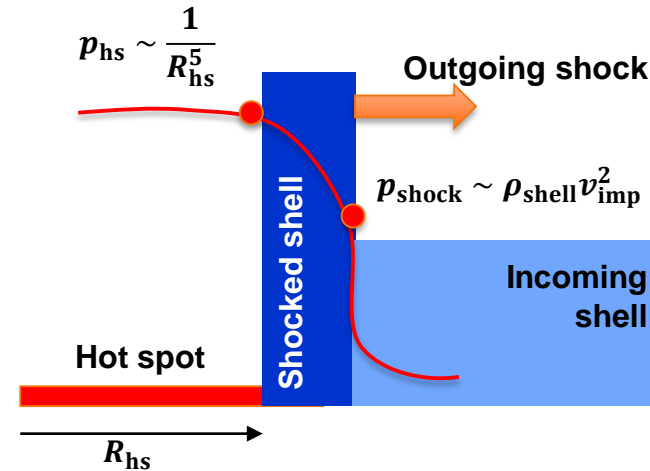
$$P_{\text{hs}} R_{\text{hs}} > 1 \text{ Gbar} \times \text{cm}, T_{\text{i}} > 4.5 \text{ keV}$$

$$E_{\text{hs}} > 16 \text{ kJ} \left(\frac{R_{\text{hs}}}{50 \mu\text{m}} \right)^2$$



- Shell must provide efficient confinement to maximize hot-spot pressure
 - Maximum hot-spot convergence must be reached before shell starts to disassemble

If the shell's dynamic pressure ($\rho_{\text{shell}} v_{\text{imp}}^2$) is too low, the shell cannot support high hot-spot pressures ($p_{\text{hs}} \propto v_{\text{imp}}^3$)



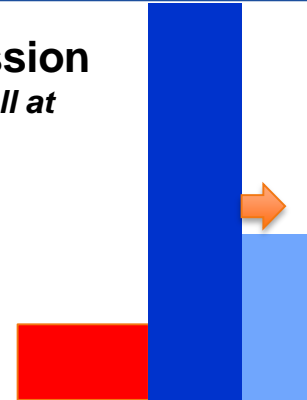
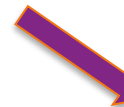
The shocked shell is decelerated by pressure gradient

$$\frac{p_{\text{hs}} - p_{\text{shock}}}{\Delta}$$

Efficient compression
Shock is inside the shell at peak compression

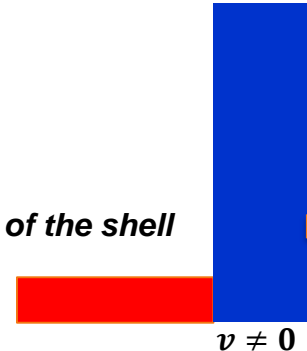


Inefficient compression
Shock breaks out of the shell before stagnation



- Maximum hot-spot pressure for given implosion parameters (v_{imp}, α)

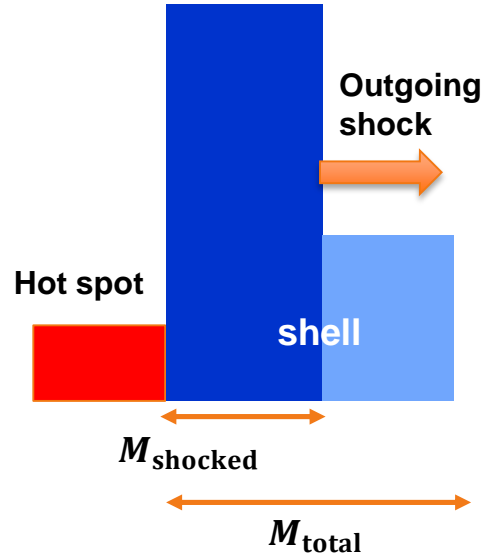
$$p_{\text{hs}} \sim \frac{p_a^{\frac{1}{3}} v_{\text{imp}}^3}{\alpha}$$



- Shell starts to disassemble from outside
- Hot-spot convergence and pressure are limited by shell mass

$$p_{\text{hs}} < \frac{p_a^{\frac{1}{3}} v_{\text{imp}}^3}{\alpha}$$

Maximum hot-spot convergence must be reached before the outgoing shock breaks out of the shell



Efficient compression:

Shock inside the shell at peak compression:

$$M_{shocked} < M_{total}$$



$$v_{imp} < 2.6 \times 10^7 \frac{\text{cm}}{\text{s}} \alpha^{0.3} \left(\frac{p_a}{100 \text{ Mbar}} \right)^{0.65}$$

Ablation pressure sets the limit on the maximum implosion velocity for an efficient piston.

What's wrong with using lower drive intensities and pressures

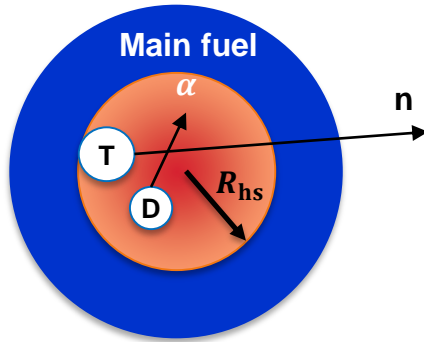
II. Creating igniting hot spot is challenging

Lower drive pressure requires **MUCH** more laser energy for ignition

$$(\rho R)_{\text{hs}} T_{\text{hs}} = 0.3 \frac{\text{g}}{\text{cm}^2} \times 5 \text{ keV} \rightarrow$$

$$P_{\text{hs}} R_{\text{hs}} > 1 \text{ Gbar} \times \text{cm}, T_{\text{i}} > 4.5 \text{ keV}$$

$$E_{\text{hs}} > 16 \text{ kJ} \left(\frac{R_{\text{hs}}}{50 \mu\text{m}} \right)^2$$



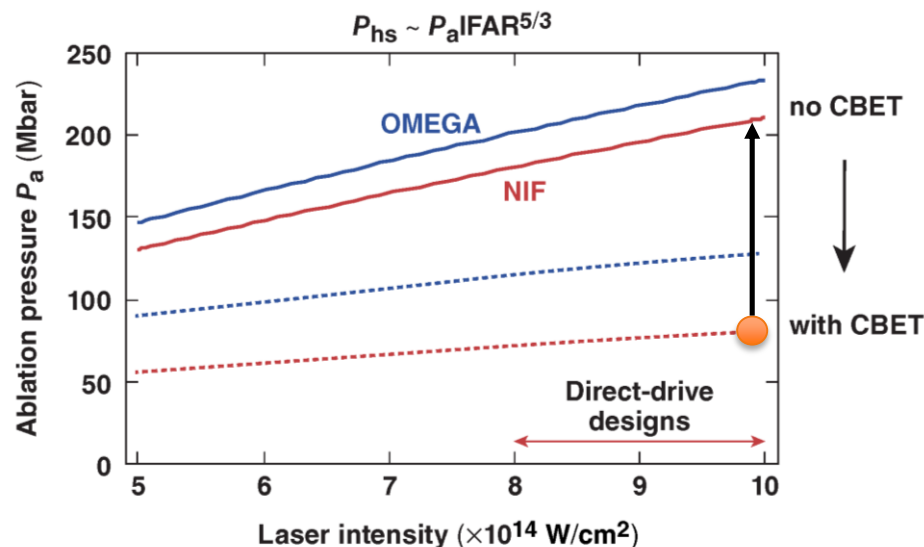
$$v_{\text{imp}} < 2.6 \times 10^7 \frac{\text{cm}}{\text{s}} \alpha^{0.3} \left(\frac{p_a}{100 \text{ Mbar}} \right)^{0.65}$$

$$E_{\text{k}}(\text{kJ}) > 156 \left(\frac{p_a}{100 \text{ Mbar}} \right)^{-5}$$

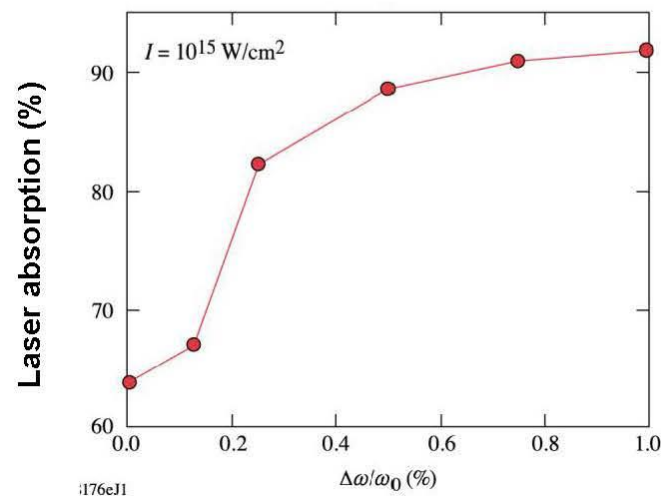
Current LID: $P_a \approx 150 \text{ Mbar} \rightarrow E_{\text{k}} > 20 \text{ kJ}$ – right at the edge

Required energy increases rapidly with reduced drive pressure $P_a = 100 \text{ Mbar} \rightarrow E_{\text{k}} > 156 \text{ kJ}$

The path forward for laser-driven IFE is to reduce LPI losses and maximize ablation pressure with broadband lasers



Crossed-beam energy transfer (Increased drive pressure)



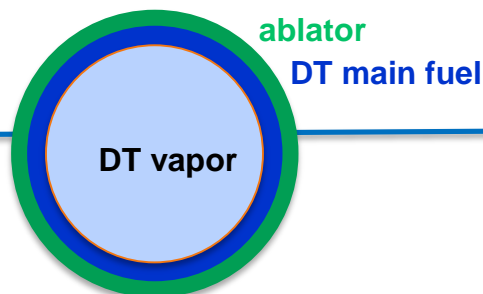
- Broadband illumination also mitigates imprint

A shorter-wavelength ArF driver also provides improved ablation pressures (see talk by S. Obenschain and whitepapers from NRL)

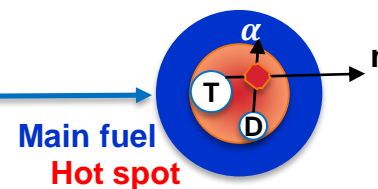
Targets for IFE must be cheap!

Hot-spot ICF approach

Conventional hot-spot ICF approach



Ignition onset



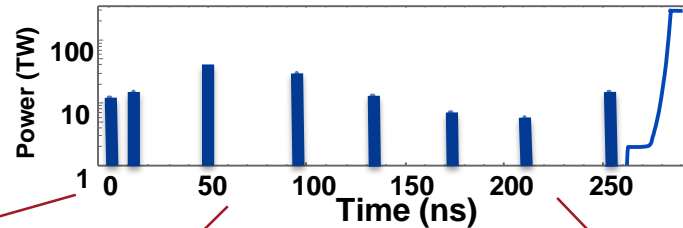
Possibly a better path for IFE: start with liquid DT ball

Wetted foam

Liquid DT

- Dynamic shell design* (ARPA-e BETHE project)
- LANL wetted-foam designs with liquid DT layers (see R. Olson white paper)

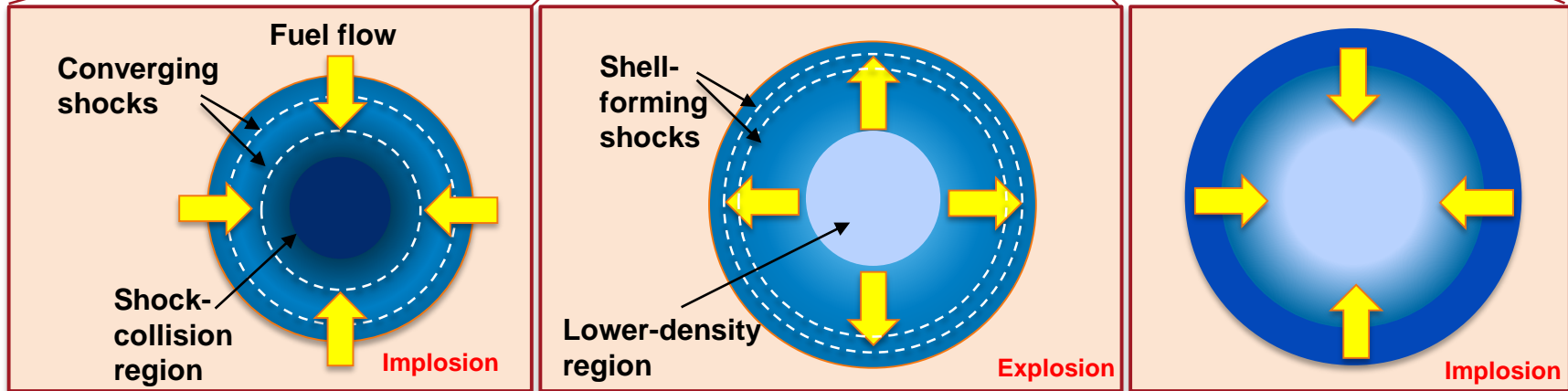
The dynamic shell design evolves through three stages



Shock heating

Blast wave expansion, density relaxation, and shell formation

Shell acceleration and hot-spot formation



Priority list for laser-drive IFE

1) Demonstrate the engine

- Broadband glass lasers (Jon Zuegel's talk)
- Excimer lasers, ArF (Steve Obenschain)

Recommendation:

- *A single-beam prototype should test feasibility of laser meeting efficiency, high rep-rate requirements*
- *A test compression facility must provide data on improved drive pressures and reduced LPI;
LPI is multibeam phenomena so single-beam test beds will not deliver the required validation*

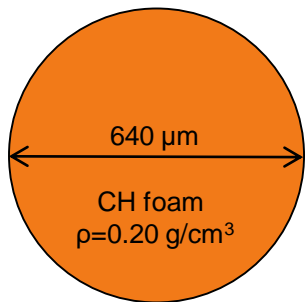
2) Demonstrate feasibility of cheap targets

- Recommendations:**
- *Wetted-foam ablaters must be validated on current facilities (OMEGA)*
 - *Several stages of advanced designs (i.e., dynamic shell concept) can be tested on current lasers (proof-of-principle experiments)*

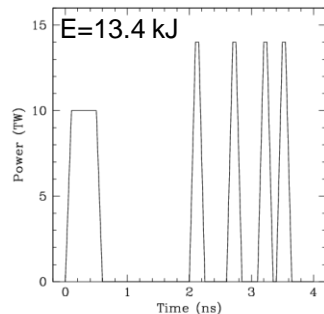
BACKUP slides

Proof-of-principle experiments are scheduled on OMEGA on 16 August 2022 as part of the LBS program

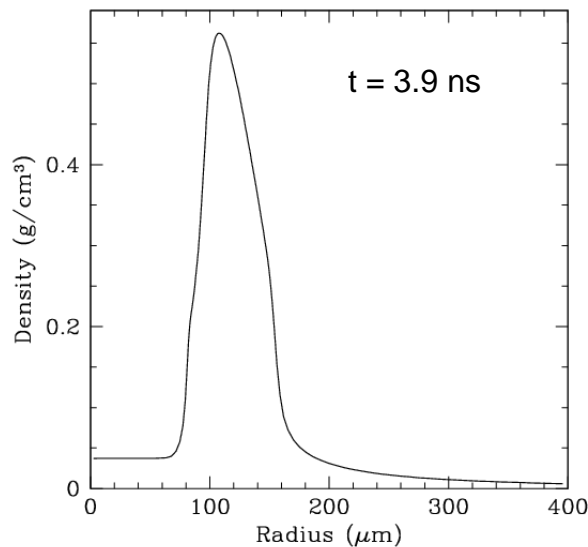
Target



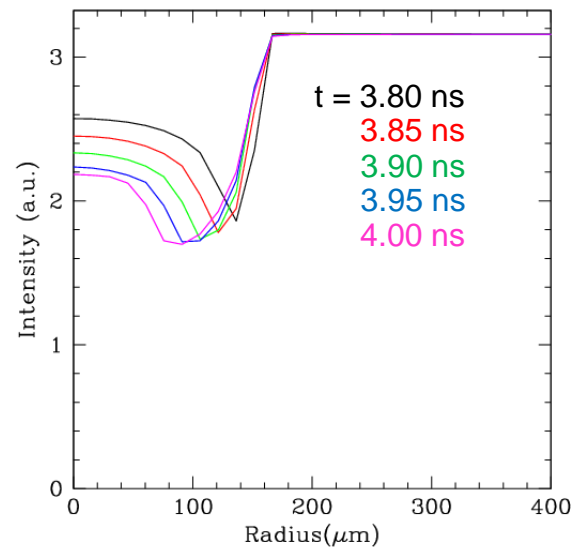
Laser pulse using SG5-650



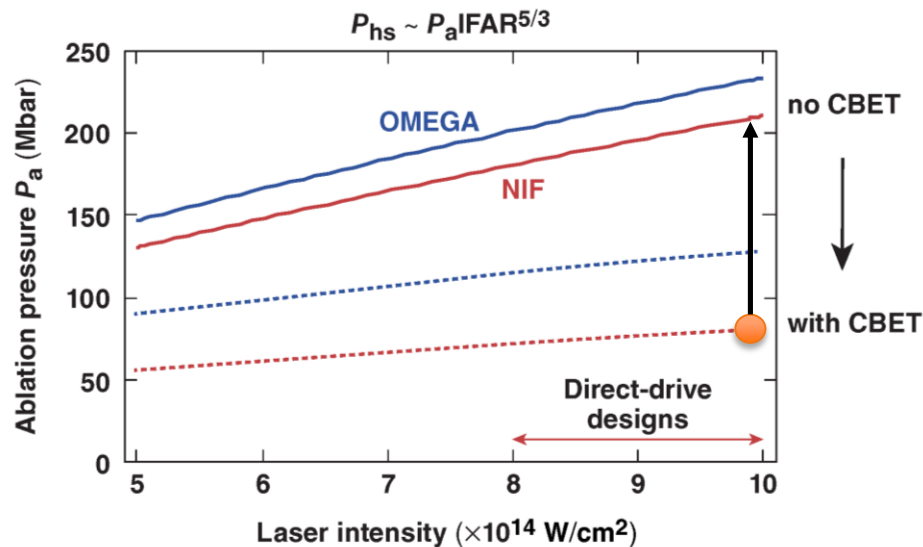
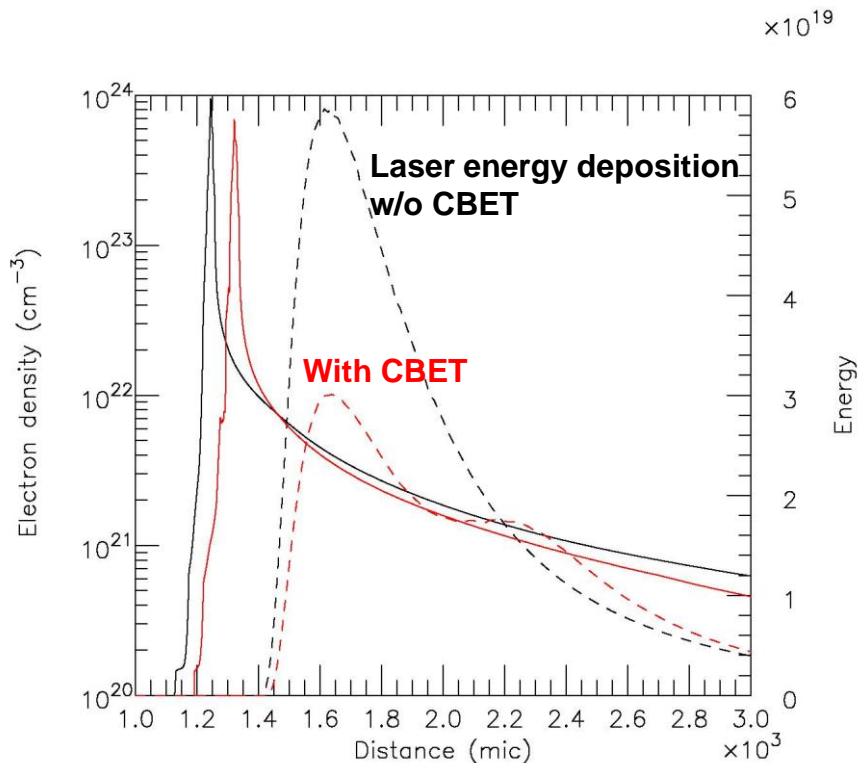
Density



Backlighting at 1.85 keV (400 eV source)



LPI coupling losses shift the laser deposition region further out into plasma corona



Further improvement in ablation pressure can be achieved through beam zooming; shorter-wavelength driver (ArF) is also helpful

Pulse shape consists of set of pickets and the main drive pulse

$$\rho R_f = 0.5 \frac{g}{cm^2}$$

$$Y_n = 6 \times 10^{14}$$

